

# Development of a Model for Restructuring of Fractal Soot Aggregates and its Parameterization using AFM Experiments

Egor Demidov<sup>\*</sup>, Ali Hasani, Alexei Khalizov, Gennady Gor

January 20, 2024



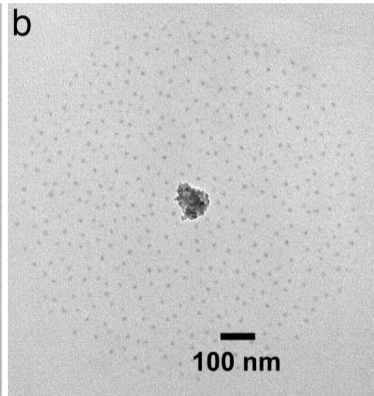
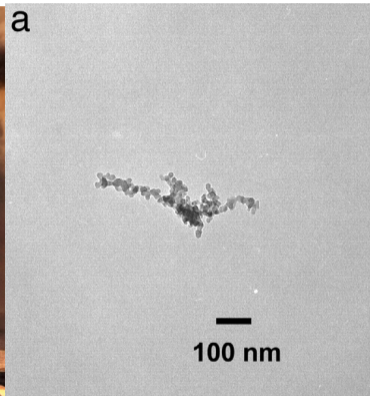
# Table of contents

- 1 Background
- 2 Discrete element method / particle dynamics
- 3 Parametrization
- 4 Supplemental
- 5 Bibliography

# Soot and its morphology

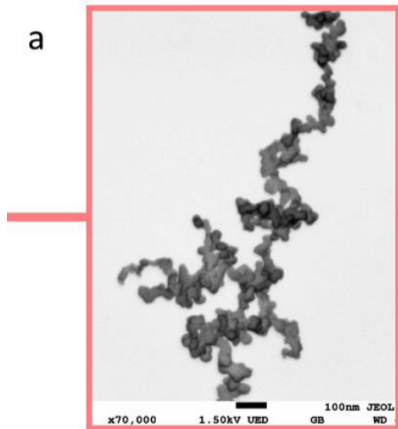


# Soot and its morphology



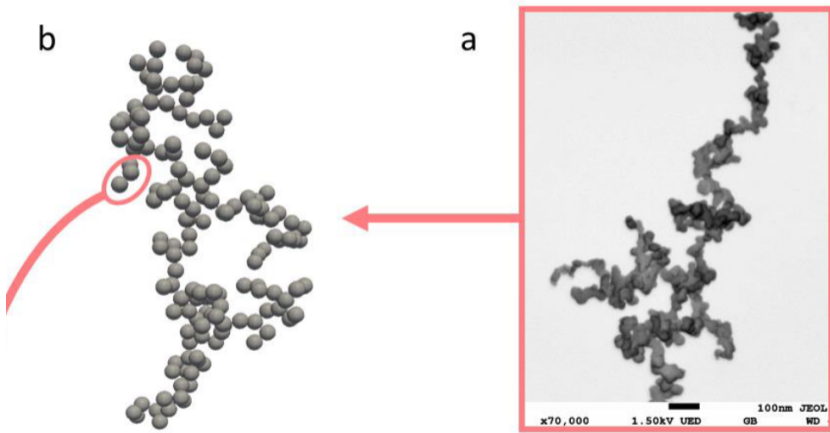
<sup>0</sup>Renyi Zhang et al. "Variability in morphology, hygroscopicity, and optical properties of soot aerosols during atmospheric processing". In: *Proceedings of the National Academy of Sciences* 105.30 (2008), pp. 10291–10296

# Soot and its morphology



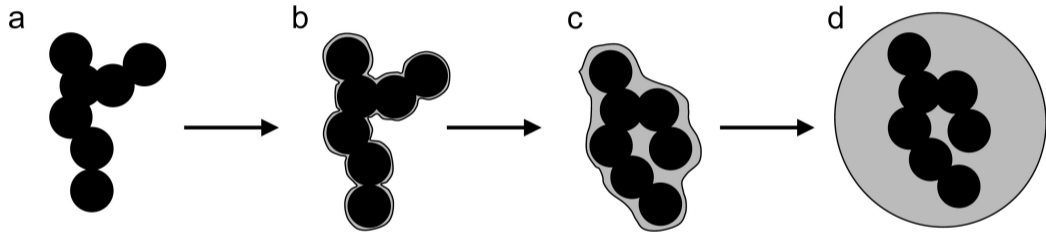
Soot is a fractal-like agglomerate of carbon particles

# Soot and its morphology



We can represent it as an agglomerate of rigid spherical particles

# Soot and its morphology



Condensation drives restructuring of agglomerates (the agglomerate rearranges due to forces exerted by the liquid)

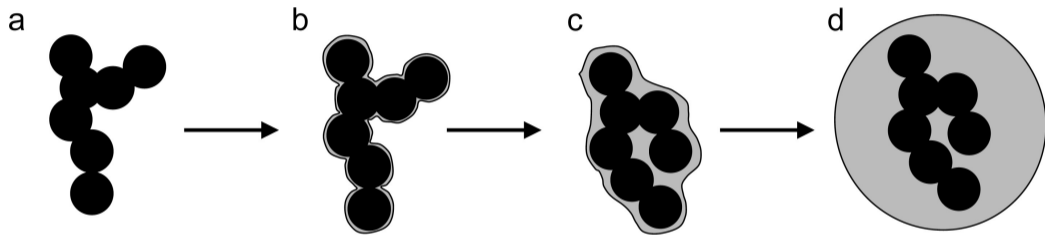
Supplemental: animations

# Table of contents

- 1 Background
- 2 Discrete element method / particle dynamics**
- 3 Parametrization
- 4 Supplemental
- 5 Bibliography



# Particle dynamics



- We model restructuring by
  - Calculating the forces acting on every particle
  - Integrating 2<sup>nd</sup> Law

$$m \frac{d^2 \mathbf{x}}{dt^2} = \sum_i \mathbf{f}_i$$

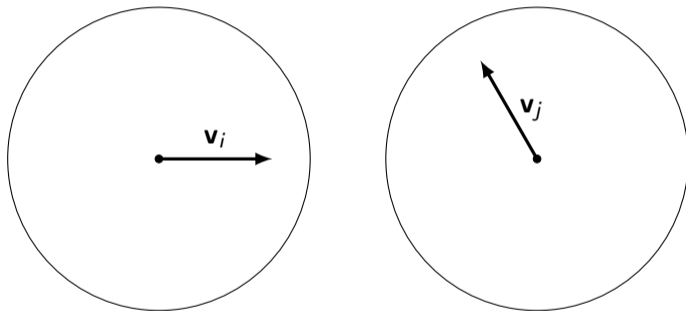
$$I \frac{d^2 \theta}{dt^2} = \sum_i \tau_i$$

$$m \frac{d^2 \mathbf{x}}{dt^2} = \sum_i \mathbf{f}_i$$

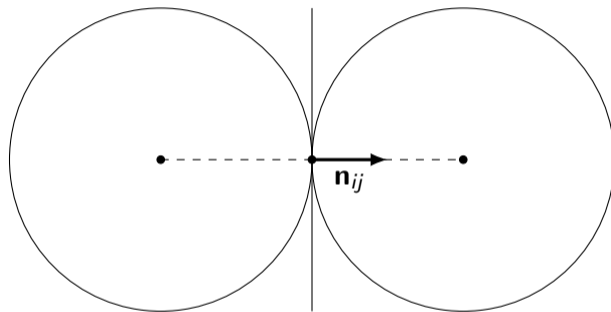
$$I \frac{d^2 \theta}{dt^2} = \sum_i \tau_i$$

- Unlike in molecular dynamics
  - Particles (atoms) are not point-masses
  - Particles have nonzero moment of inertia and undergo rotational motion
  - Frictional forces between contacting particles need to be considered

# Contact forces

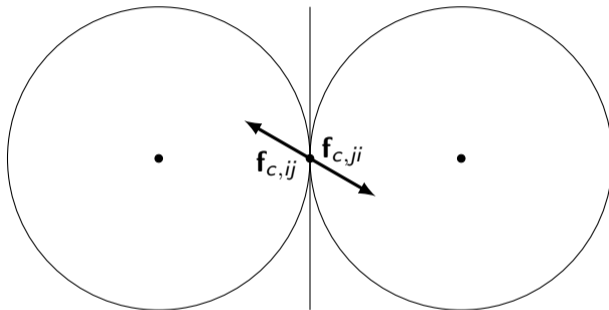


## Contact forces



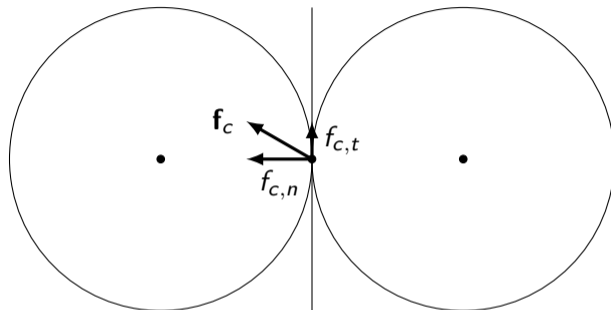
Contact plane between particles  $i$  and  $j$  is defined by normal vector  $\mathbf{n}_{ij}$  connecting centers of respective particles

## Contact forces



Contact will result in particle  $i$  experiencing force  $\mathbf{f}_{c,ij}$  and particle  $j$  experiencing an equal in magnitude and opposite force  $\mathbf{f}_{c,ji} = -\mathbf{f}_{c,ij}$  applied at the contact point

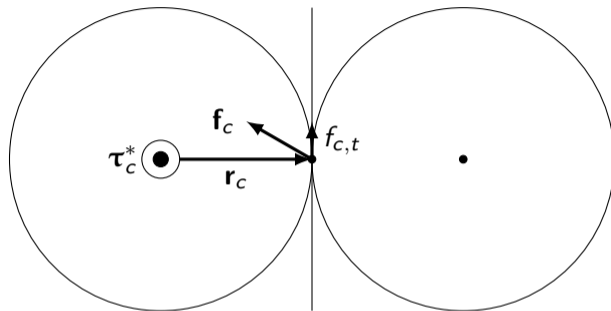
## Contact forces



Contact force  $\mathbf{f}_c$  can be decomposed into normal and tangential components  $f_{c,n}$  and  $f_{c,t}$ , where  $f_{c,n}$  is related to elasticity and  $f_{c,t}$  is related to friction:

$$\mathbf{f}_c = f_{c,n}\mathbf{n} + f_{c,t}\mathbf{t}$$

## Contact forces



\*  $\tau_c$  is directed out-of-plane

Since the contact force is applied to the contact point, it will result in a torque:

$$\tau_c = \mathbf{r}_c \times \mathbf{f}_c = f_{c,t} (\mathbf{r}_c \times \mathbf{t})$$



# Contact forces

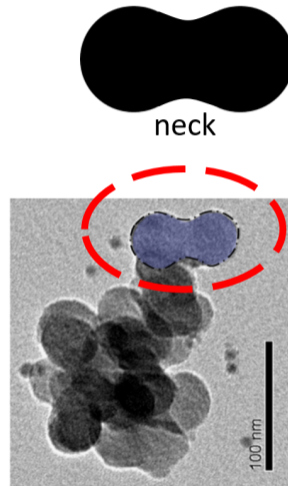
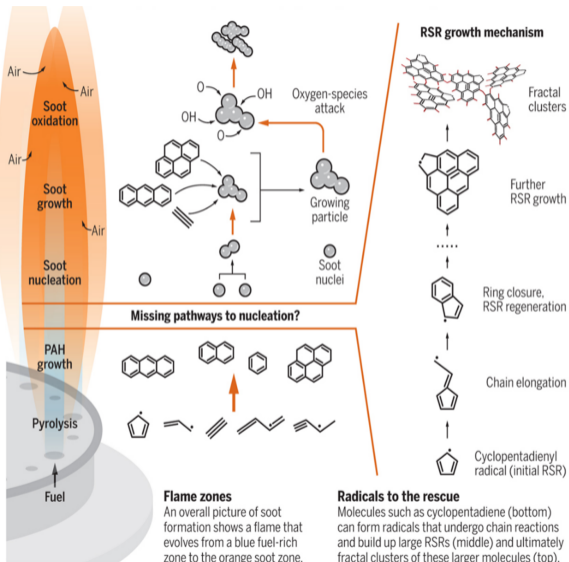
Recap:

- Contact forces are expressed in terms of elasticity and friction coefficients
- They only act on free particles that come into contact
- Various models exist, including ones that simulate energy dissipation

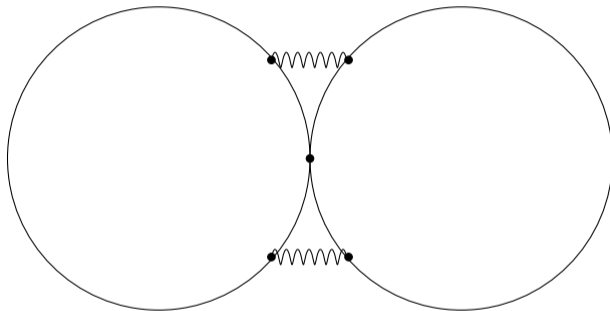
[Supplemental: animations](#)

[Supplemental: Luding 2008](#)

# Sinter bridges

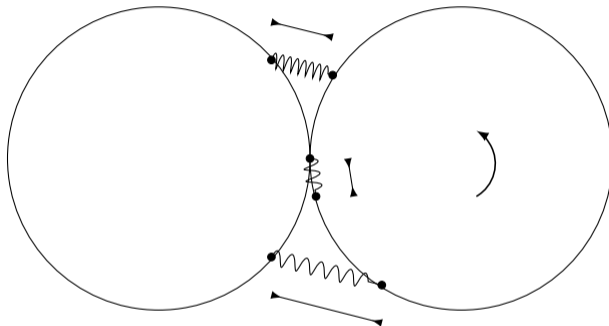


# Sinter bridges



Three (2D) / five (3D) springs are inserted between sintered particles

## Sinter bridges



Translation / rotation of one particle gives rise to forces that restore the reference configuration of the pair

# Sinter bridges

## Numerical implementation:

Motion of a spring connector  $\mathbf{s}$  attached to particle  $i$  with center at  $\mathbf{x}$ :

$$\frac{d\mathbf{s}}{dt} = \mathbf{v}_i + \boldsymbol{\omega}_i \times (\mathbf{s} - \mathbf{x}_i)$$

Spring displacement, force, and torque:

$$\Delta l = \|\mathbf{s}_i - \mathbf{s}_j\| - l_0$$

$$\mathbf{f}_n = -k_n \Delta l \frac{\mathbf{s}_i - \mathbf{s}_j}{\|\mathbf{s}_i - \mathbf{s}_j\|}$$

$$\boldsymbol{\tau}_n = \mathbf{r}_n \times \mathbf{f}_n$$

# Sinter bridges

Breaking criterion:

Potential energy stored in a neck - sum of PEs stored in constituent springs:

$$E_{\text{neck}} = \sum E_{\text{spring}}$$

$$E_{\text{spring}} = \frac{1}{2} k_n (l - l_0)^2$$

Let us break the neck when

$$E_{\text{neck}} \geq E_{\text{crit}}$$

where  $E_{\text{crit}}$  is an attribute of the neck and is assigned at its creation

Supplemental: animations

## Non-contact / body forces

Hamaker (Van der Waals) attraction:

$$\mathbf{f}_{VdW} = \frac{2AR}{(\|\mathbf{x}_i - \mathbf{x}_j\| - 2R)^2} \mathbf{n}$$

if  $\|\mathbf{x}_i - \mathbf{x}_j\| - 2R > h_0$

$$\mathbf{f}_{VdW} = \frac{2AR}{h_0^2} \mathbf{n}$$

if  $\|\mathbf{x}_i - \mathbf{x}_j\| - 2R \leq h_0$

Coating forces: to be implemented

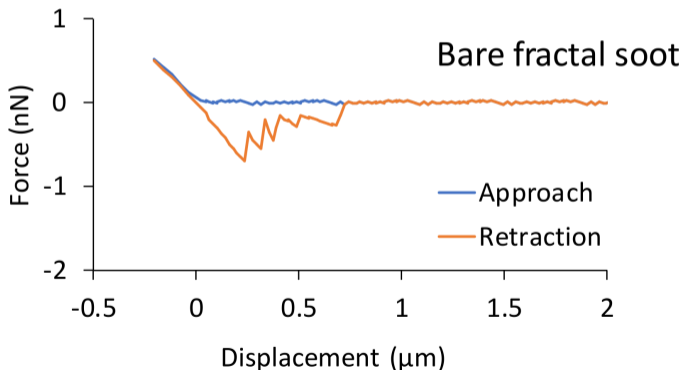
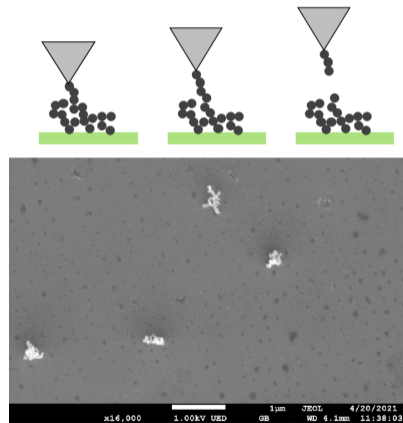
# Table of contents

- 1 Background
- 2 Discrete element method / particle dynamics
- 3 Parametrization**
- 4 Supplemental
- 5 Bibliography



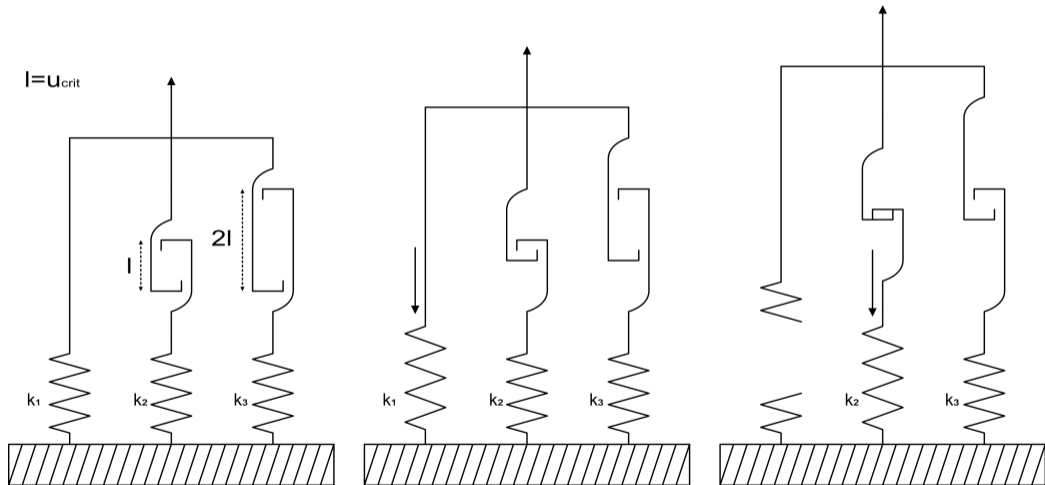
## Mechanics only simulations

Before coating is implemented, AFM experiments<sup>1</sup> and TEM tomography can be used to parametrize the model:

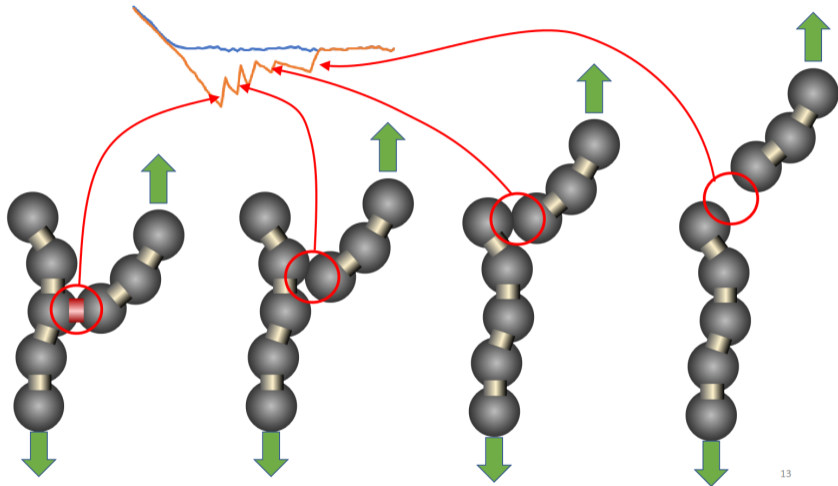


<sup>1</sup>Khalizov and Hasani, AAAR 2023

# Mechanics only simulations



# Mechanics only simulations

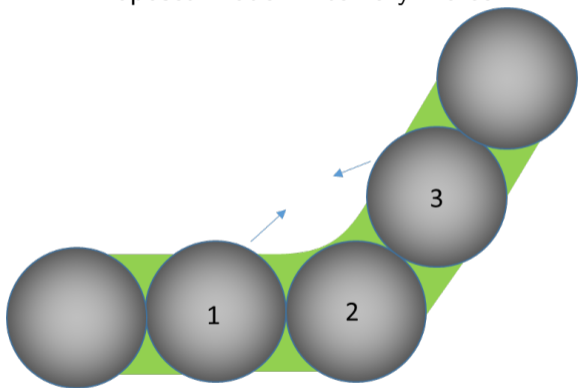


13

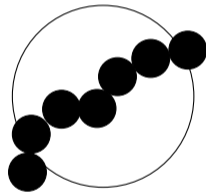
Supplemental: animations

# Adding capillary forces to the DEM model

Proposed model - "ternary" force:



Proposed model - interface force:



# Table of contents

- 1 Background
- 2 Discrete element method / particle dynamics
- 3 Parametrization
- 4 Supplemental**
- 5 Bibliography

# Animations

- Contact forces (Luding 2008)
  - [No resistance](#)
  - [Rolling resistance](#)
  - [Torsion resistance](#)
- Necks
  - [Trimer](#)
  - [Y-shape](#)
  - [Aggregate \(central force\)](#)

Return: soot and its morphology

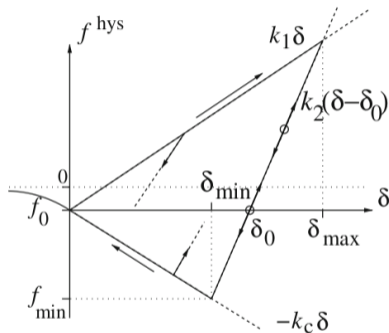
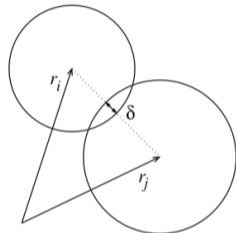
Return: contact forces

Return: sinter bridges

Return: mechanics only simulations

# Luding 2008<sup>1</sup> contact model

Normal force, friction, rolling and torsion resistance



Return: contact forces

<sup>1</sup>Stefan Luding. "Cohesive, frictional powders: contact models for tension". In: *Granular matter* 10.4 (2008), pp. 235–246

# AFM transfer function

Cantilever dynamics:

$$m \frac{d^2 x}{dt^2} + \gamma \frac{dx}{dt} + kx = g(t)$$

where  $x$  is cantilever displacement and  $g(t)$  is the forcing function on the cantilever.

Apparent force  $f(t)$ :

$$f(t) = kx(t)$$

$$\frac{m}{k} \frac{d^2 f}{dt^2} + \frac{\gamma}{k} \frac{df}{dt} + f(t) = g(t)$$



# AFM transfer function

For a critically damped cantilever:

$$\gamma = 2\sqrt{mk}$$

Transfer function becomes:

$$\frac{F(s)}{G(s)} = \frac{\omega_0^2}{(s + \omega_0)^2}$$

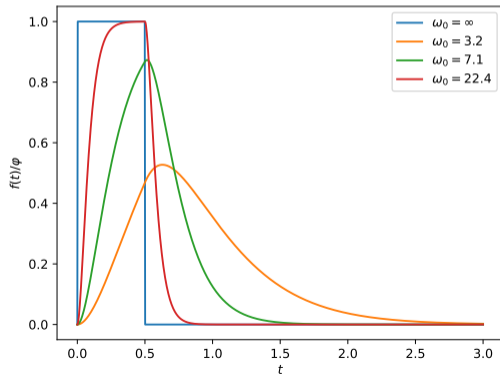
Where  $\omega_0$  is the natural angular frequency of the undamped system:

$$\omega_0 = \sqrt{\frac{k}{m}}$$

# AFM transfer function

Response to step input (Heaviside step function):

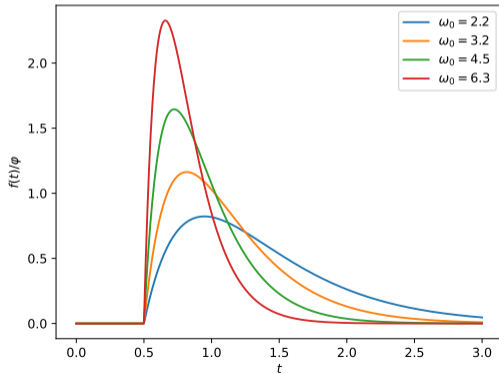
$$f(t) = \varphi (1 - e^{-\omega_0 t} (\omega_0 t + 1))$$



# AFM transfer function

Response to unit impulse (Dirac delta function):

$$f(t) = \varphi\omega_0^2 e^{-\omega_0 t}$$



# Table of contents

- 1 Background
- 2 Discrete element method / particle dynamics
- 3 Parametrization
- 4 Supplemental
- 5 Bibliography**

# Bibliography I



Luding, Stefan. “Cohesive, frictional powders: contact models for tension”. In: *Granular matter* 10.4 (2008), pp. 235–246.



Zhang, Renyi et al. “Variability in morphology, hygroscopicity, and optical properties of soot aerosols during atmospheric processing”. In: *Proceedings of the National Academy of Sciences* 105.30 (2008), pp. 10291–10296.